# THE g-RECURRENT CONDITION IMPOSED ON THE EINSTEIN'S CONNECTION

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#### 1. Introduction.

Various recurrent connections have been studied by many authors, such as Chung, Datta, E.M. Patterson, M. Prvanovitch, Singal, and Takano, etc(refer to [4] and [5]). Examples of such connections are that of Ricci-recurrent curvature, that of birecurrent curvature, and skew-symmetric recurrent connection. In this paper, we introduce a new concept of g-recurrent connection in a generalized n-dimensional Riemannian manifold  $X_n$ , and prove that g-recurrent condition imposed on the Einstein's connection is meaningless from the physical point of view.

#### 2. Preliminaries.

This section is a brief collection of definitions and notations which are needed in our subsequent considerations. Let  $X_n$  be a generalized n-dimensional Riemannian manifold referred to a real coordinate system  $x^{\nu}$ , which obeys only coordinate transformations  $x^{\nu} \to x^{\nu'}$  for which

$$(2.1)  $Det(\frac{\partial x'}{\partial x}) \neq 0$$$

The manifold  $X_n$  is endowed with a general real nonsymmetric tensor  $g_{\lambda\mu}$  which may be split into its symmetric part  $h_{\lambda\mu}$  and skew-symmetric part  $k_{\lambda\mu}^{-1}$ 

$$(2.2) g_{\lambda\mu} = h_{\lambda\mu} + k_{\lambda\mu}$$

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<sup>&</sup>lt;sup>1</sup>Throughout the present paper, all Greek indices take the values 1, 2, ..., n and follow the summation convention unless stated otherwise.

where

(2.3) 
$$g = Det(g_{\lambda\mu}) \neq 0, \quad \mathfrak{h} = Det(h_{\lambda\mu}) \neq 0, \quad \mathfrak{k} = Det(k_{\lambda\mu})$$

In virtue of (2.3) we may define a unique tensor  $h^{\lambda\nu}$  by

$$(2.4) h_{\lambda\mu}h^{\lambda\nu} = \delta^{\nu}_{\mu}$$

which together with  $h_{\lambda\mu}$  will serve for raising and/or lowering indices of tensors in X in the usual manner. There exists also a unique tensor  $*q^{\lambda\nu}$  satisfying

$$(2.5) g_{\lambda\mu}^{\phantom{\lambda\mu}*}g^{\lambda\nu} = g_{\mu\lambda}^{\phantom{\mu\lambda}*}g^{\nu\lambda} = \delta^{\nu}_{\mu}$$

The manifold  $X_n$  is connected by a general real connection  $\Gamma_{\lambda}{}^{\nu}{}_{\mu}$  with the following transformation rule:

(2.6) 
$$\Gamma_{\lambda'}{}^{\nu'}{}_{\mu'} = \frac{\partial x^{\nu'}}{\partial x^{\alpha}} \left( \frac{\partial x^{\beta}}{\partial x^{\lambda'}} \frac{\partial x^{\gamma}}{\partial x^{\mu'}} \Gamma_{\beta}{}^{\alpha}{}_{\gamma} + \frac{\partial^{2} x^{\alpha}}{\partial x^{\lambda'} \partial x^{\mu'}} \right)$$

It may also be decomposed into its symmetric part  $\Lambda_{\lambda}{}^{\nu}{}_{\mu}$  and its skew-symmetric part  $S_{\lambda\mu}{}^{\nu}$ , called the torsion tensor of  $\Gamma_{\lambda}{}^{\nu}{}_{\mu}$ :

$$(2.7) \qquad \Gamma_{\lambda}{}^{\nu}{}_{\mu} = \Lambda_{\lambda}{}^{\nu}{}_{\mu} + S_{\lambda\mu}{}^{\nu}; \quad \Lambda_{\lambda}{}^{\nu}{}_{\mu} = \Gamma_{(\lambda}{}^{\nu}{}_{\mu)}; \quad S_{\lambda\mu}{}^{\nu} = \Gamma_{[\lambda}{}^{\nu}{}_{\mu]}$$

A connection  $\Gamma_{\lambda}{}^{\nu}{}_{\mu}$  is said to be *einstein* if it satisfies the following system of Einstein's equations:

(2.8a) 
$$\partial_{\omega}g_{\lambda\mu} - \Gamma_{\lambda}{}^{\alpha}{}_{\omega}g_{\alpha\mu} - \Gamma_{\omega}{}^{\alpha}{}_{\mu}g_{\lambda\alpha} = 0$$

or equivalently

$$(2.8b) D_{\omega}g_{\lambda\mu} = 2S_{\omega\mu}{}^{\alpha}g_{\lambda\alpha}$$

where  $D_{\omega}$  is the symbolic vector of the covariant derivative with respect to  $\Gamma_{\lambda}{}^{\nu}{}_{\mu}$ . The manifold  $X_n$  connected by this Einstein's connection is a generalization of the space-time  $X_4$ , and Einstein's n-dimensional unified field theory is based upon this manifold  $X_n$ . Our new concept of g-recurrent connection  $\Gamma_{\lambda}{}^{\nu}{}_{\mu}$  is defined by the following system of equations:

$$(2.9) D_{\omega}g_{\lambda\mu} = 2X_{\omega}g_{\lambda\mu}$$

for a non-null vector  $X_{\mu}$ . The manifold  $X_n$  connected by this connection is called an *n*-dimensional *g-recurrent manifold*.

The main purpose of the present paper is to prove that the Einstein's connection satisfying the g-recurrent condition (2.9) is meaningless from the physical point of view.

# 3. The g-recurrent connection.

This section is devoted to the investigations of the differential geometric properties of g-recurrent connections. The following two theorems will be proved simultaneously:

THEOREM 3.1. The system (2.9) may be decomposed into

$$(3.1a) D_{\omega}h_{\lambda\mu} = 2X_{\omega}h_{\lambda\mu}$$

(3.1b) 
$$D_{\omega}k_{\lambda\mu} = 2X_{\omega}k_{\lambda\mu}.$$

THEOREM 3.2. The system (2.9) is equivalent to

$$(3.2) D_{\omega}^* g^{\lambda \nu} = -2X_{\omega}^* g^{\lambda \nu}.$$

*Proof.* The equations (3.1a,b) follow from (2.9) and

$$D_{\omega}h_{\lambda\mu} = D_{\omega}g_{(\lambda\mu)}, \quad D_{\omega}k_{\lambda\mu} = D_{\omega}g_{[\lambda\mu]}.$$

In virtue of (2.5), multiplication of  ${}^*g^{\lambda\nu}$  to both sides of (2.9) gives

$$(3.3) -g_{\lambda\mu}D_{\omega}^*g^{\lambda\nu} = {}^*g^{\lambda\nu}D_{\omega}g_{\lambda\mu} = 2X_{\omega}g_{\lambda\mu}^*g^{\lambda\nu} = 2X_{\omega}\delta_{\mu}^{\nu}$$

The equations (3.2) may be obtained by multiplying  ${}^*g^{\epsilon\mu}$  again to both sides of (3.3). Conversely, start with (3.2), and multiply this equations by  $g_{\lambda\mu}$  to get (2.9).

REMARK 3.3. The form of equations (3.2) may be used for the study of g-recurrent connections in the Einstein's n-dimensional g-unified field theory(Refer to [1],[2],[10]).

The following scalars will be used in our subsequent considerations:

$$(3.4) g = \frac{\mathfrak{g}}{\mathfrak{h}}, \quad k = \frac{\mathfrak{k}}{\mathfrak{h}}$$

THEOREM 3.4. The covariant derivative of the determinants g and h are

$$(3.5a) D_{\omega}\mathfrak{g} = 2n\mathfrak{g}X_{\omega}$$

$$(3.5b) D_{\omega}\mathfrak{h} = 2n\mathfrak{h}X_{\omega}.$$

*Proof.* We first note that a direct consequence of (2.9) is

(3.6) 
$$D_{\omega}\mathfrak{g} = \frac{\partial \mathfrak{g}}{\partial q_{\lambda\mu}} D_{\omega} g_{\lambda\mu} = \mathfrak{g}^* g^{\lambda\mu} D_{\omega} g_{\lambda\mu}$$

On the other hand, multiplication of  ${}^*g^{\lambda\mu}$  to both sides of (2.9) gives

$$(3.7) *g^{\lambda\mu} D_{\omega} g_{\lambda\mu} = 2n X_{\omega}$$

The relation (3.5a) immediately follows by substituting (3.7) into (3.6). The relation (3.5b) may be proved similarly by starting from (2.4) and (3.1)a.

THEOREM 3.5. If the system (2.9) admits a solution  $\Gamma_{\lambda}{}^{\nu}{}_{\mu}$ , it must be of the form

(3.8) 
$$\Gamma_{\lambda \mu}^{\nu} = \{\lambda_{\mu}^{\nu}\} + S_{\lambda \mu}^{\nu} + V_{\lambda \mu}^{\nu}$$

where  $\{\lambda^{\nu}_{\mu}\}$  are the Christoffel symbols with respect to  $h_{\lambda\mu}$  and

(3.9) 
$$V^{\nu}{}_{\lambda\mu} = V^{\nu}{}_{(\lambda\mu)} = -2S^{\nu}{}_{(\lambda\mu)} - 2X_{(\lambda}\delta_{\mu)}{}^{\nu} + X^{\nu}h_{\lambda\mu}$$

Proof. In virtue of

$$D_{\omega}h_{\lambda\mu} = \partial_{\omega}h_{\lambda\mu} - \Gamma_{\lambda}{}^{\alpha}{}_{\omega}h_{\alpha\mu} - \Gamma_{\mu}{}^{\alpha}{}_{\omega}h_{\lambda\alpha}$$

We have

$$\frac{1}{2}h^{\nu\alpha}(D_{\lambda}h_{\alpha\mu} + D_{\mu}h_{\lambda\alpha} - D_{\alpha}h_{\lambda\mu})$$

$$= \{\lambda^{\nu}_{\mu}\} - 2h^{\nu\alpha}S_{\alpha(\lambda\mu)}\Gamma_{(\lambda^{\nu}_{\mu})}$$

$$= \{\lambda^{\nu}_{\mu}\} - 2S^{\nu}_{(\lambda\mu)} - \Gamma_{\lambda^{\nu}_{\mu}} + S_{\lambda\mu}^{\nu}$$
(3.10)

On the other hand, the relation (3.1a) gives

(3.11) 
$$\frac{1}{2}h^{\nu\alpha}(D_{\lambda}h_{\alpha\mu} + D_{\mu}h_{\lambda\alpha} - D_{\alpha}h_{\lambda\mu}) = 2X_{(\lambda}\delta_{\mu)}^{\ \nu} - X^{\nu}h_{\lambda\mu}$$

Comparing (3.10) and (3.11), we finally have (3.8) in virtue of (3.9).

REMARK 3.6. In virtue of (3.8) and (3.9), we note that the investigation of the g-recurrent connection  $\Gamma_{\lambda}{}^{\nu}{}_{\mu}$  is reduced to the study of the tensor  $S_{\lambda\mu}{}^{\nu}$ . In order to know the g-recurrent connection  $\Gamma_{\lambda}{}^{\nu}{}_{\mu}$ , it is necessary and sufficient to represent the tensor  $S_{\lambda\mu}{}^{\nu}$  in terms of  $g_{\lambda\mu}$ . This is an open problem. Probably, the precise tensorial representation of  $S_{\lambda\mu}{}^{\nu}$  in terms of  $g_{\lambda\mu}$  may be obtained by starting from (3.1b).

# 4. The g-recurrent condition imposed on the Einstein's connection.

In this section, we investigate the meaning of g-recurrent condition in the Einstein's n-dimensional unified field theory from physical point of view.

THEOREM 4.1. A necessary condition for the system (2.9) to admit a solution is that the scalars g and k, defined by (3.4), are constant.

*Proof.* In virtue of Theorem 3.1, we note that the system (2.9) is equivalent to (3.1a,b). Multiplication of  ${}^*g^{\lambda\mu}$  to both sides of (2.9) gives

$$2nX_{\omega} = (\partial_{\omega}g_{\lambda\mu} - \Gamma_{\lambda}{}^{\alpha}{}_{\omega}g_{\alpha\mu} - \Gamma_{\mu}{}^{\alpha}{}_{\omega}g_{\lambda\alpha})^{*}g^{\lambda\mu}$$

$$= (\partial_{\omega}g_{\lambda\mu})^{*}g^{\lambda\mu} - 2\Gamma_{\alpha}{}^{\alpha}{}_{\omega}$$

$$= \frac{1}{\mathfrak{g}}(\partial_{\omega}g_{\lambda\mu})\frac{\partial\mathfrak{g}}{\partial g_{\lambda\mu}} - 2\Gamma_{\alpha}{}^{\alpha}{}_{\omega}$$

$$= \partial_{\omega}(\ln\mathfrak{g}) - 2\Gamma_{\alpha}{}^{\alpha}{}_{\omega}$$

$$(4.1)$$

Similarly, multiplying  $h^{\lambda\mu}$  to both sides of (3.1a), we have

$$(4.2) 2nX_{\omega} = \partial_{\omega}(\ln \mathfrak{h}) - 2\Gamma_{\alpha}{}^{\alpha}{}_{\omega}$$

Comparing (4.1) and (4.2), we have

(4.3) 
$$\partial_{\omega}(\ln \mathfrak{g}) = \partial_{\omega}(\ln \mathfrak{h}) \text{ or } q = \text{constant}$$

which proves the first statement. If k=0, then our theorem is proved. If  $k\neq 0$ , then there exists a unique inverse tensor  $\bar{k}^{\lambda\mu}$  such that

$$(4.4) k_{\lambda\alpha}\bar{k}^{\nu\alpha} = \delta^{\nu}_{\lambda}$$

Consequently, multiplying  $\bar{k}^{\lambda\mu}$  to both sides of (3.1b) it follows that (4.5)  $2nX_{\omega} = \partial_{\omega}(\ln \mathfrak{k}) - 2\Gamma_{\alpha}{}^{\alpha}{}_{\omega}$ 

which together with (4.2) give

k = constant.

REMARK 4.2. In the Einstein's unified field theory, a function of scalar g may be identified with the gravitational function (Refer to [7],[9]). Therefore, if we assume that Einstein's connection is also grecurrent in the Einstein's unified field theory, the gravitational function is reduced to a constant in the gravitational theory in virtue of Theorem 4.1. From the physical point of view, this is a strong restriction to the generality of Einstein's unified field theory. Consequently, the adoptation of the condition (2.9) in the Einstein's unified field theory is meaningless.

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