# ON THE KNOTTED ELASTIC CURVES 

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#### Abstract

According to the Bernoulli-Euler theory of elastic rods the bending energy of the wire is proportional to the total squared curvature of $\gamma$, which we will denote by $F(\gamma)=\int_{\gamma} k^{2} d s$. If the result of J.Langer and D.Singer [3] extend to knotted elastic curve, then we obtain the following. Let $\{\gamma, M\}$ be a closed knotted elastic curve. If the curvature of $\gamma$ is nonzero for everywhere, then $\gamma$ lies on torus.


## I. Introduction

Elastic curve (or elastica) and its generalizations have long been of interest in the context of elasticity theory. The elastica as a purely geometrical entity seems to have been largely ignored (for historical references concerning the classical elastica, we refer to the recent survey by Truesdel [2]).

Elastic curve is a mathematical model of Peano curve. And elastic energy (bending energy) is critical for $\mathcal{T}$ defined on regular curves. Euler was able to obtain a good qualitative description of all plane elastic curves. In fact, Peano curve not only has a curve but also knot.Thus elastic curve is not complete mathematical model of Peano curve. Here, in order to establish a mathematical model, consider the energy which is the sum of elastic energy and knotted energy. And define the curve its energy is critical.

## II. Main theorem

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All curves, functions, vecterfields will be assumed $C^{\infty}$ class. For 3 -dimensional Euclidean space $\mathbf{R}^{3}$, Euclidean inner product will be denoted by $<,>$ and the Euclidean norm by $|\mid$.

Let $\gamma=\gamma(t):\left[t_{1}, t_{2}\right] \rightarrow \mathbf{R}^{3}$ be a holomorphic $C^{\infty}$-class curve. $v$ will denote the velocity of $\gamma$ and $T, k$ will denote unite tangent vector and curvature of $\gamma$. (i.e, $v=\left|\frac{d \gamma}{d t}\right|, T=\frac{1}{v} \frac{d \gamma}{d t}, k=\left|\nabla_{T} T\right|$ where $\nabla_{T}=\frac{1}{v} \frac{d}{d t}$ )

Define functional $\mathcal{F}$ by $\mathcal{F}(\gamma)=\int_{t_{1}}^{t_{2}} k^{2} v d t$ and we called it elastic energy of $\gamma$.

Let $M$ be a unit normal vector field along $\gamma$ and $\{\gamma, M\}$ be a curve with unit normal vector field.

Definition 1. Let $\{\gamma, M\}$ be a curve with unit normal vectorfield and its domain is $\left[t_{1}, t_{2}\right]$. Define a function $h$ on $\left[t_{1}, t_{2}\right]$ by $h=<$ $\nabla_{T} M, L>$, and we called it a knot function of $\{\gamma, M\}$. Here $L=$ $T \times M, \times$ is exterior product in $\mathbf{R}^{3} . h(t)$ be a quantity of knot of $M$ at $\gamma(t)$.

Remark. If $M$ is parallel to normal connection along to $\gamma$, then $h \equiv 0$.

Definition 2. $\{\gamma, M\}, v, h$ are the same notation as above.
(1) $\int_{t_{1}}^{t_{2}} h^{2} v d t$ is called a knotted energy of $\{\gamma, M\}$.
(2) Let $\epsilon>0$ be a constant. Define a functional $\mathcal{T}_{\epsilon}$ with respect to curve with unit normal vector field by

$$
\mathcal{T}_{\epsilon}(\{\gamma, M\})=\mathcal{T}(\gamma)+\epsilon \int_{t_{1}}^{t_{2}} h^{2} v d t
$$

$\mathcal{T}_{\epsilon}(\{\gamma, M\})$ is called knotted elastic energy of coefficient $\epsilon$ of $\{\gamma, M\}$. Here, domain of $\mathcal{T}_{\epsilon}$ is the set of all curve with unit normal vector field.

Definition 3. Let $t_{0}>0, \phi \in \mathbf{R} / 2 \pi \mathbf{Z}$. $\{\gamma, M\}$ is called period $t_{0}$ if the following two conditions are satisfied.
(1) $\gamma\left(t+t_{0}\right)=\gamma(t)$.
(2) $M\left(t+t_{0}\right)=R(\phi) M(t)$ where $R(\phi): T^{2} \mathbf{R}^{3} \rightarrow T^{\perp} \mathbf{R}^{3}$ be a rotation of angle $\phi$ in each fiber of normal bundle $T^{\perp} \mathbf{R}^{3}$ along $\gamma$. (orientation is $\left.R\left(\frac{\pi}{2}\right) M(t)=L(t)\right)$

Let $C\left(t_{0}, \phi\right)$ be a set of all curves with unit normal vector field its period is $t_{0}$. Then we obtain a following lemma from the a first variation formula.

Lemma 1. $\{\gamma, M\} \in C\left(t_{0}, \phi\right)$ (its length $\int_{0}^{t_{0}} v d t$ is fixed) is critical point of $\mathcal{T}_{\epsilon}$ iff there exist real numbers $\mu, \sigma$ such that the following are satisfied.
(1) $\nabla_{T}\left[2\left(\nabla_{T}\right)^{2} T+\left(3 k^{2}-\mu+\epsilon h^{2}\right) T-2 \epsilon h R\left(\frac{\pi}{2}\right)\left(\nabla_{T} T\right)\right]=0$.
(2) $h(t)=\sigma$.

Definition 4. Let $\{\gamma, M\}$ be a curve with unit normal vector field with velocity 1 (i.e. $v \equiv 1$ ). If (1) and (2) of the Lemma 1 are satisfied, then $\{\gamma, M\}$ is called a knotted elastic curve, $\sigma$ is called a knot parameter of $\{\gamma, M\}$.

Let $l>0, \phi \in \mathbf{R} / 2 \pi \mathbf{Z}$ and $\mathcal{U} C(l, \phi)$ be the set of velocity 1 of $\gamma$ in $C(l, \phi)$ and $\{\gamma, M\}$ be an element of $\mathcal{U} C(l, \phi)$.

Define

$$
\begin{gathered}
T_{\{\gamma, M\}} C(l, \phi)=\left\{\begin{array}{l}
\Lambda \text { is vector field of period } \\
(\Lambda, f) \mid l \text { along } \gamma, f \text { is a function } \\
\text { of period } l
\end{array}\right\}, \\
T_{\{\gamma, M\}} \mathcal{U} C(l, \phi)=\left\{(\Lambda, f) \left\lvert\, \begin{array}{l}
(\Lambda, f) \in T_{\{\gamma, M\} C(l, \phi)} \\
\text { i.e. }<\nabla_{T} \Lambda, T>=0
\end{array}\right.\right\} .
\end{gathered}
$$

Then the following lemma is satisfied.
Lemma 2. For the variation $\{\gamma, M\}_{\lambda}$ of $\{\gamma, M\}$ in $\mathcal{U} C(l, \phi)$, $\{\gamma, M\}_{\lambda}=\left\{\gamma_{\lambda}, M_{\lambda}\right\}\left(-\lambda_{0}<\lambda<\lambda_{0},\{\gamma, M\}_{0}=\{\gamma, M\}\right)$
$\left(\left.\frac{\partial \gamma_{l}}{\partial \lambda}\right|_{\lambda=0},<\frac{\partial M_{\lambda}}{\partial \lambda}, L_{\lambda}>\left.\right|_{\lambda=0}\right) \in T_{\{\gamma, M\}} \mathcal{U} C(l, \phi)$. Left side is called a variational vector field of variation $\{\gamma, M\}_{\lambda}$

Conversely, for any $(\Lambda, f) \in T_{\{\gamma, M\}} \mathcal{U C}(l, \phi)$, there exist a variation $\{\gamma, M\}_{\lambda}$ of $\{\gamma, M\}$ in $\mathcal{U} C(l, \phi)$ such that $\left.\frac{\partial \gamma_{l}}{\partial \lambda}\right|_{\lambda=0}=\Lambda,\left\langle\frac{\partial M_{\lambda}}{\partial \lambda}, L_{\lambda}\right\rangle$ $\left.\right|_{\lambda=0}=f$.

In the above Lemma 2, $T_{\{\gamma, M\}} \mathcal{U} C(l, \phi)$ is tangent space of $\mathcal{U C}(l, \phi)$ at $\{\gamma, M\}$.

Lemma 3. Let $\{\gamma, M\}$ be a knotted elastic curve and also variational vector field of $\{\gamma, M\}_{\lambda}$. Then

$$
\left.\frac{d^{2}}{d \lambda^{2}}\right|_{\lambda=0} \mathcal{T}_{\epsilon}\left(\{\gamma, M\}_{\lambda}\right)=\int_{0}^{l}<\mathcal{T}_{\{\gamma, M\}}(\Lambda, f),(\Lambda, f)>d s
$$

where

$$
\begin{aligned}
& \mathcal{T}_{\{\gamma, M\}}(\Lambda, f)=\left(p \left[\nabla _ { T } \left\{2\left(\nabla_{T}\right)^{3} \Lambda+\left(3 k^{2}-\mu+\epsilon \sigma^{2}\right) \nabla_{T} \Lambda\right.\right.\right. \\
&-2 \epsilon \sigma R\left(\frac{\pi}{2}\right)\left(\left(\nabla_{T}\right)^{2} \Lambda-<\left(\nabla_{T}\right)^{2} \Lambda, T>T\right) \\
&\left.\left.-2 \epsilon\left(<\nabla_{T} \Lambda, R\left(\frac{\pi}{2}\right)\left(\nabla_{T} T\right)>+T f\right) R\left(\frac{\pi}{2}\right)\left(\nabla_{T} T\right)\right\}\right] \\
&\left.-2 \epsilon\left(T<\nabla_{T} \Lambda, R\left(\frac{\pi}{2}\right)\left(\nabla_{T} T\right)>+T^{2} f\right)\right) \\
&\left(p: T_{\{\gamma, M\}} C(l, \phi) \rightarrow T_{\{\gamma, M\}} \mathcal{U} C(l, \phi)\right. \text { is an orthogonal projection }
\end{aligned}
$$ with respect to $L^{2}$-inner product)

We consider eigenvalue problem $\mathcal{T}_{\{\gamma, M\}}(\Lambda, f)=p(\Lambda, f), \phi \in \mathbf{R}$, we can obtain the following theorem.

Theorem 4. Let $\gamma$ be a circle with radius 1 and $\{\gamma, M\}$ be a knotted elastic curve. Then the eigenvalue of Jacobi operator $\mathcal{I}_{\{\gamma, M\}}$ has the following properties.
(1) The eigenvalue of $\mathcal{T}_{\{\gamma, M\}}$ is positive whenever $0 \leq \epsilon^{2} \sigma^{2}<3$.
(2) There exists a non-trivial eigenvector its eigen value 0 whenever $\epsilon^{2} \sigma^{2}=m^{2}-1(2 \leq m, m$ is integer $)$.
(3) There exists a negative eigenvalue whenever $\epsilon^{2} \sigma^{2}>3$.

The curvature $k$ and torsion $\tau$ of elastic curve are represented by elliptic function. By J. Langer and D. Singer [3], for every elastic curve $\gamma$ in $\mathbf{R}^{3}$ there is naturally associated to $\gamma$ a cylindrical coordinate system $(r, \theta, z)$ on $\mathbf{R}^{3}$, the restrictions to $\gamma$ of the coordinate fields $\frac{\partial}{\partial r}, \frac{\partial}{\partial \theta} \frac{\partial}{\partial z}$ being expressible in terms of $k, \tau, T, N, B$. Thus we can see the following theorem.

Theorem 5 [J. Langer and D. Singer][3]. Let $\gamma$ be a closed elastic curve. Then $\gamma$ lies on embedded tori of revolution.

If the proof of the above theorem extend to a knotted elastic curve, then we obtain the following result.

Theorem 6. Let $\{\gamma, M\}$ be a closed knotted elastic curve. Suppose that, the curvature $k$ of $\gamma$ is nonzero for everywhere. Then $\gamma$ lies on torus.

Proof. Let $\{\gamma, M\}$ be a knotted elastic curve. First of all, find the curvature $k$ and torsion $\tau$. Application of the Frenet formula for (1) of [Lemma 1] leads to the ordinary differential equation for $k, \tau$. Solution of the equation is concretely represented by Jacobi function. In this representation, $k$ and $\tau$ are periodic functions with the same period and if $k$ is constant then $\tau$ is also constant. Secondly, we will construct a cylindrical coordinate system. If $k$ is constant, then $\gamma$ is a straight line or a circle. Suppose that $k$ is not constant.

If $\gamma=\gamma(s)$ be a curve with velocity 1 and its curvature is positive at every point, Then $\Lambda$ is a vector field along $\gamma . \Lambda$ extends to a killing vector field on $\mathbf{R}^{3}$ iff $\Lambda$ satisfies the following.

$$
\begin{gather*}
<\nabla_{T} \Lambda, T>=0  \tag{a}\\
<\left(\nabla_{T}\right)^{2} \Lambda, N>=0  \tag{b}\\
<\left(\nabla_{T}\right)^{3} \Lambda-\frac{k_{s}}{k}\left(\nabla_{T}\right)^{2} \Lambda+k^{2} \nabla_{T} \Lambda, B>=0 \tag{c}
\end{gather*}
$$

where $T, N, B$ form the Frenet frame for $\gamma$. Put

$$
\begin{gathered}
J_{0}=2\left(\nabla_{T}\right)^{2} T+\left(3 k^{2}-\mu+\epsilon \sigma^{2}\right) T-2 \epsilon \sigma R\left(\frac{\pi}{2}\right)\left(\nabla_{T} T\right) \\
H=\epsilon \sigma T+k B \\
J_{1}=H-\frac{<J_{0}, H>}{\left|J_{0}\right|^{2}} J_{0}(\mu, \sigma \text { are constant in Lema } 1) .
\end{gathered}
$$

Then $J_{0}, H, J_{1}$ is also killing along $\gamma$. Let $\bar{J}_{0}, \bar{J}_{1}$ be the extension of $J_{0}, J_{1}$ on $\mathbf{R}^{3}$. In (1) of Lemma 1 , we can see $\bar{J}_{0}$ is constant vector field. Thus $\bar{J}_{1}$ is a rotation field perpendicular to $\bar{J}_{0}$. By the
above statement we obtain a cylindrical coordinate $(r, \theta, z)$. It satisfies $\frac{\partial}{\partial z}=\frac{1}{\left|\bar{J}_{0}\right|} \bar{J}_{0}, \frac{\partial}{\partial \theta}=c \bar{J}_{1}$ where $c$ is a positive constant. Setting $\gamma(s)=(r(s), \theta(s), z(s))$ one then obtains

$$
\begin{gathered}
r(s)=c\left|J_{1}(s)\right| \\
\theta_{s}(s)=\frac{<T, \frac{\partial}{\partial \theta}>}{\left|\frac{\partial}{\partial \theta}\right|^{2}}=\frac{<T, J_{1}(s)>}{c\left|J_{1}(s)\right|^{2}} \\
z_{s}(s)=<T, \frac{1}{\left|J_{0}(s)\right|} J_{0}(s)>
\end{gathered}
$$

where $J_{0}, J_{1}$ are vector fields along $\gamma$ and components of $T, N, B$ are represented by curvature and torsion of $\gamma$.

If $\{\gamma, M\}$ is periodic, then $\mathrm{r}, \mathrm{z}$ are also periodic and curve of $\gamma$ in $r z-$ plane is a simple closed curve. Thus every closed knotted elastic curve lies on torus of revolution. Since curvature and torsion of knotted elastic curve are periodic function with some period, $\gamma$ is periodic iff $\Delta z=0$ (i.e. $\frac{\Delta \theta}{2 \pi} \in \mathbf{Q}$ )

Theorem 7 [J.Langer and D.Singer][3]. For any closed elastic curve $-\pi \leq \Delta \theta \leq 0, \frac{\Delta \theta}{2 \pi} \in \mathbf{Q}$ and conversely for any $\psi$, such that $-\pi \leq \psi \leq 0, \frac{\psi}{2 \pi} \in \mathbf{Q}$ there exists a unique closed elastic curve such that $\Delta \theta=\psi$.

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